

# New results and open problems toward a minimal realization of the light composite Higgs

Lattice Higgs Collaboration (LatHC)

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**Lattice for Beyond the Standard Model Physics** 

April 23-25, 2015, LLNL workshop

the Higgs doublet field

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} \pi_2 + i \, \pi_1 \\ \sigma - i \, \pi_3 \end{pmatrix} \qquad \frac{1}{\sqrt{2}} \left( \sigma + i \, \vec{\tau} \cdot \vec{\pi} \right) \equiv M$$

$$D_{\mu}M = \partial_{\mu}M - igW_{\mu}M + ig'MB_{\mu}$$
, with  $W_{\mu} = W_{\mu}^{a}\frac{\tau^{a}}{2}$ ,  $B_{\mu} = B_{\mu}\frac{\tau^{3}}{2}$ 

The Higgs Lagrangian is

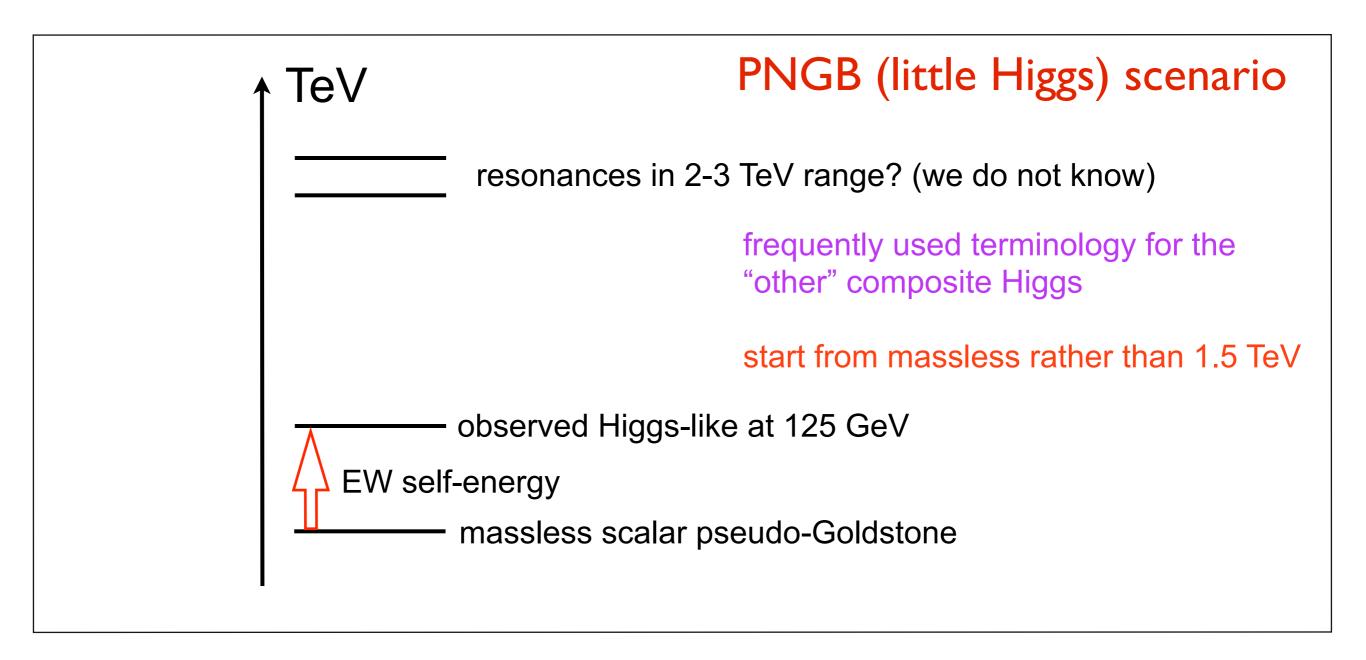
### spontaneous symmetry breaking

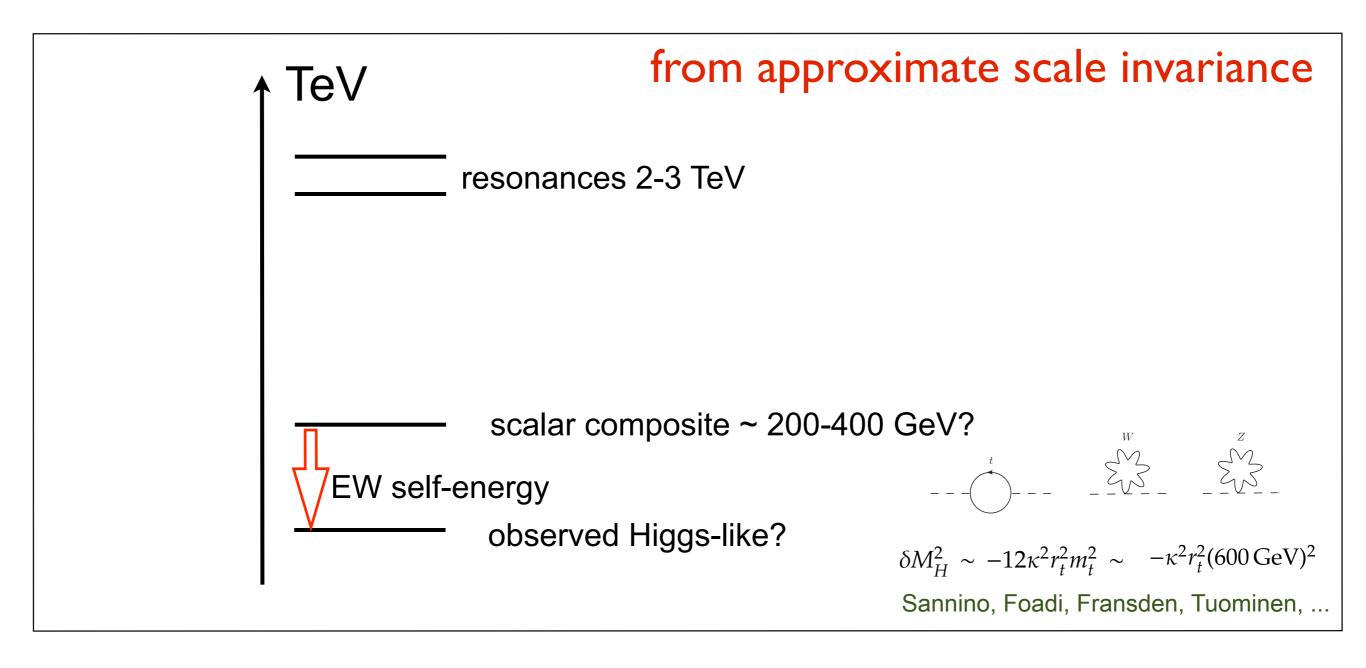
Higgs mechanism

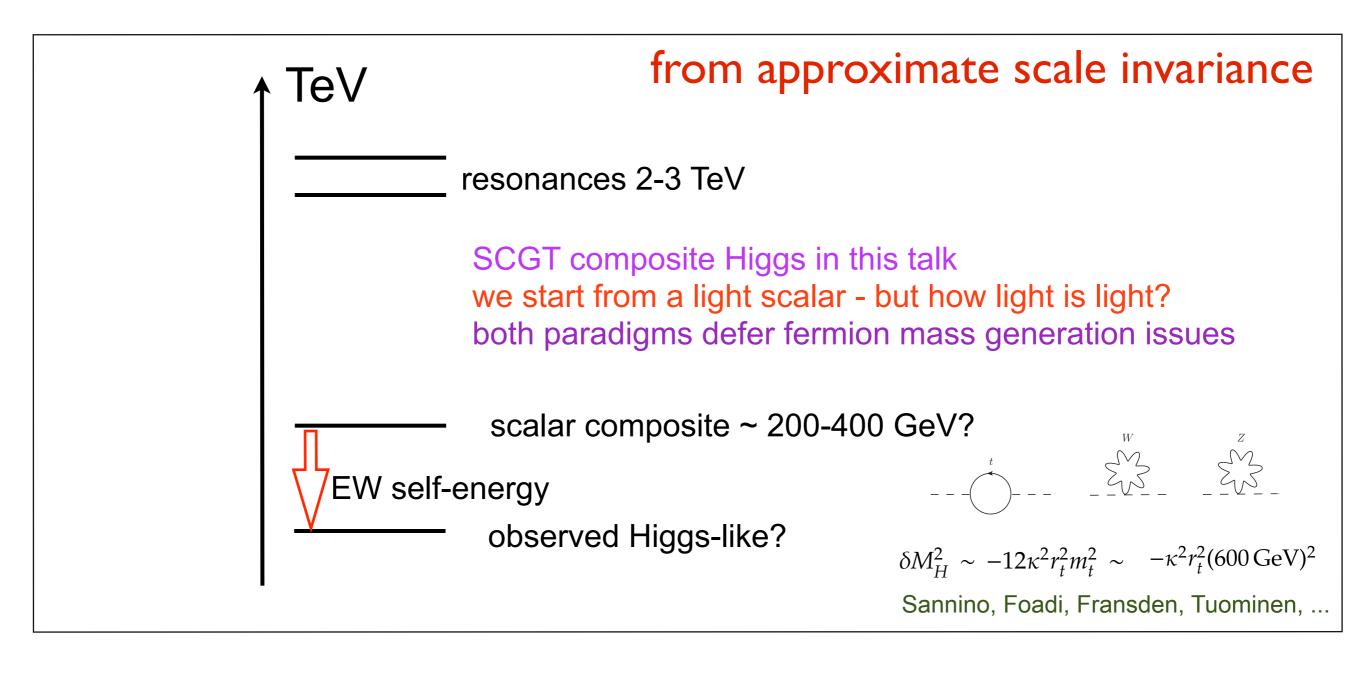
$$\mathcal{L} = \frac{1}{2} \text{Tr} \left[ D_{\mu} M^{\dagger} D^{\mu} M \right] - \frac{m_M^2}{2} \text{Tr} \left[ M^{\dagger} M \right] - \frac{\lambda}{4} \text{Tr} \left[ M^{\dagger} M \right]^2$$

strongly coupled gauge theory 
$$\mathcal{L}_{Higgs} \rightarrow -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i \bar{Q} \gamma_{\mu} D^{\mu} Q + \dots$$

needle in the haystack? or, just one of the haystacks?







### **Near-conformal SCGT?**

- light scalar close to conformal window? the D-word
- navigating mine fields of p, ε, and δ regimes in chiPT
- scale setting and spectroscopy
- mixed action strategy

### Chiral Higgs condensate

- new stochastic method for spectral density
- GMOR and mode number
- epsilon regime and RMT
- large mass anomalous dimension?

### Running coupling

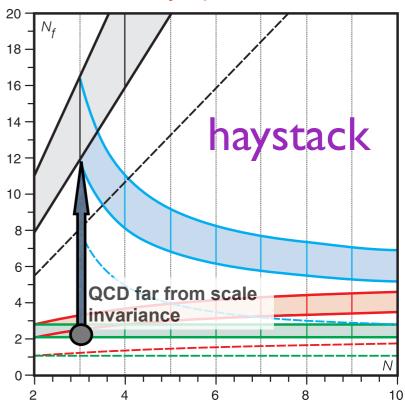
- scale dependent running coupling the R-word
- matching with mass anomalous dimension?

### Early universe

- sextet EW phase transition
- sextet baryon and dark matter

## Summary and Outlook

#### SCGT Theory Space © Sannino



close to scale invariance?

 $\begin{array}{l} \text{nf=2 sextet rep} \\ \text{massless fermions} \\ \text{SU(2) doublet} \end{array} \begin{bmatrix} u(+e/2) \\ d(-e/2) \end{bmatrix} \\ \begin{array}{l} \text{minimal EW} \\ \text{embedding} \end{array}$ 

3 Goldstones morph into weak bosons minimal realization

QCD intuition for near-conformal compositeness is just plain wrong

Technicolor thought to be scaled up QCD theme of the talk:

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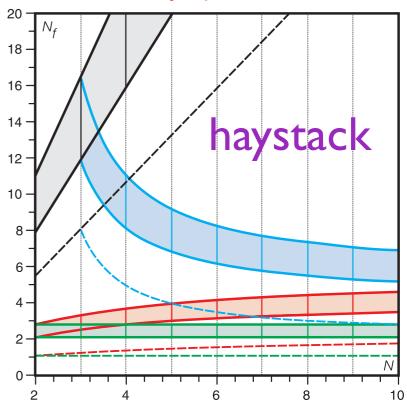
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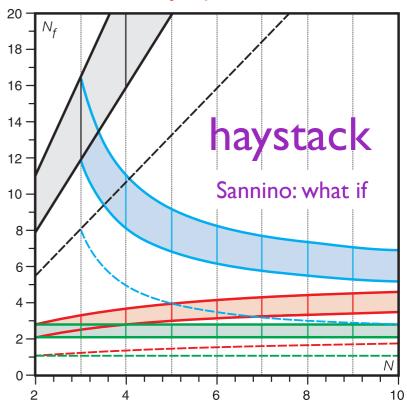
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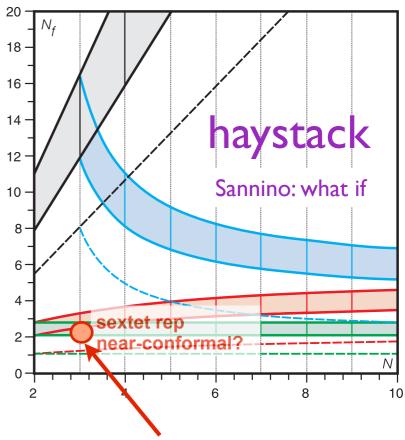
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## The light 0++ scalar

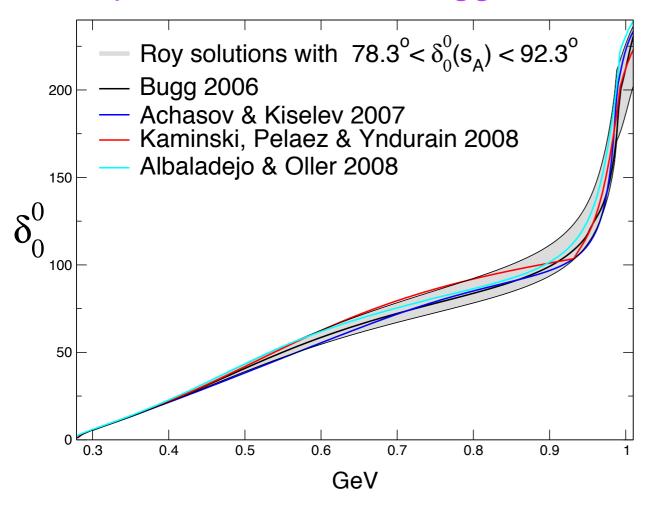
## QCD (aka old TC) 80ies,90ies

the failure of old Higgs-less technicolor: 0<sup>++</sup> scalar in QCD (bad Higgs impostor)

$$\sqrt{s_{\sigma}}$$
 = (400 - 1200) - i (250 - 500) MeV

estimate in Particle Data Book

### π-π phase shift in 0<sup>++</sup> "Higgs" channel



$$\sqrt{s_{\sigma}} = 441^{+16}_{-8} - i \, 272^{+9}_{-12.5} \,\text{MeV}$$

Leutwyler: dispersion theory combined with ChiPT

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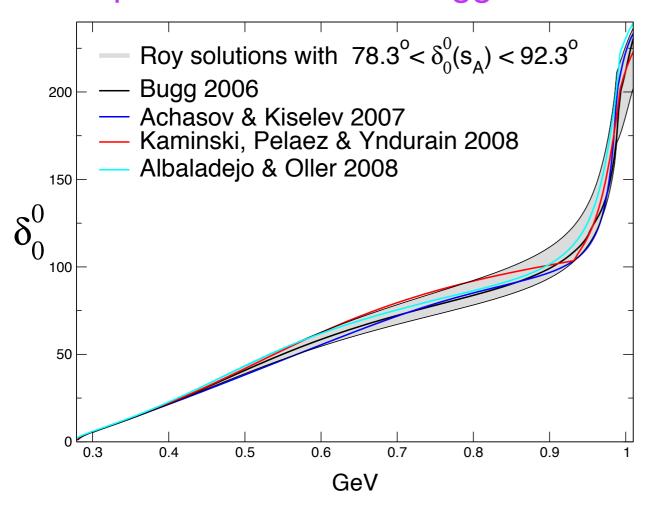
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 $\sqrt{s_{\sigma}}$  = (400 - 1200) - i (250 - 500) MeV



broad  $M_{\sigma} \sim 1.5$  TeV in old technicolor, based on scaled up QCD, hence the tag "Higgs-less"

This is expected to be different in near-conformal strongly coupled gauge theories

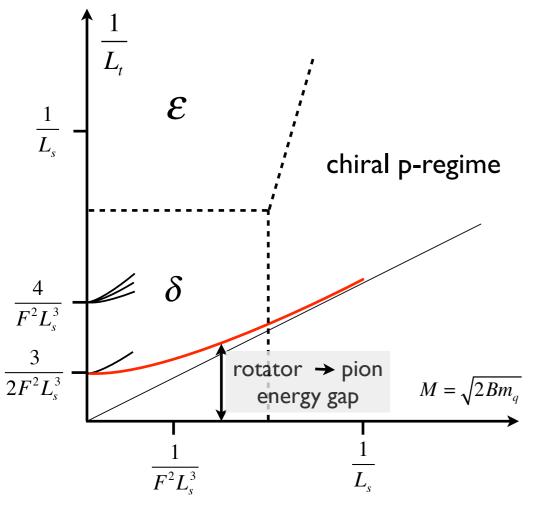
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Leutwyler: dispersion theory combined with ChiPT

## Spoiler alert:

## $FL < \sqrt{N_f/2}$ simulations $\Rightarrow$ no theory

## when in finite volume, it is always an expansion in I/FL!



Condition of reaching the chiral expansion regime can be estimated from rotator spectrum  $\Rightarrow$ 

$$E_l = \frac{1}{2\theta}l(l+2)$$
 with  $l = 0,1,2,...$  rotator spectrum for  $SU(2)_f \times SU(2)_f$ 

direct application to sextet model

$$\theta = F^2 L_s^3 (1 + \frac{C(N_f = 2)}{F^2 L_s^2} + O(1/F^4 L_s^4))$$
 (P. Hasenfratz and F. Niedermayer)

expansion in  $1/F^2L_s^2$ !

 $C(N_f = 2) = 0.45$  (FL=1 is ~ 2fm in lite QCD) C will grow with ~  $N_f$ 

the constraints are the same in the arepsilon-regime and p-regime

FL = 0.1 L=0.2 fm in QCD femto world OK to study volume dependent PT coupling running with V

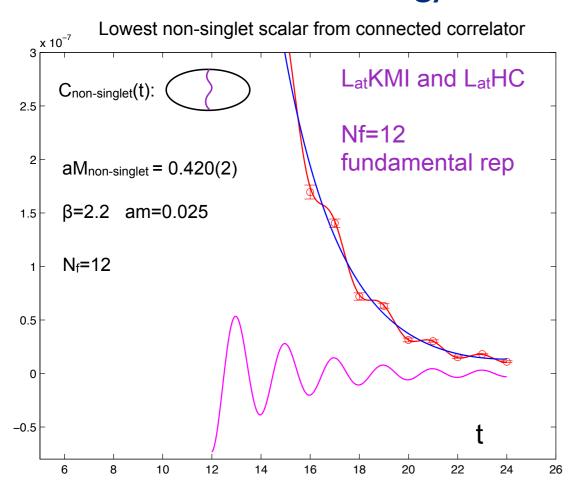
FL = I L= 2 fm in QCD and we crossed over to the  $\chi$ SB phase all 3 regimes ( $\epsilon, \delta, p$ ) OK

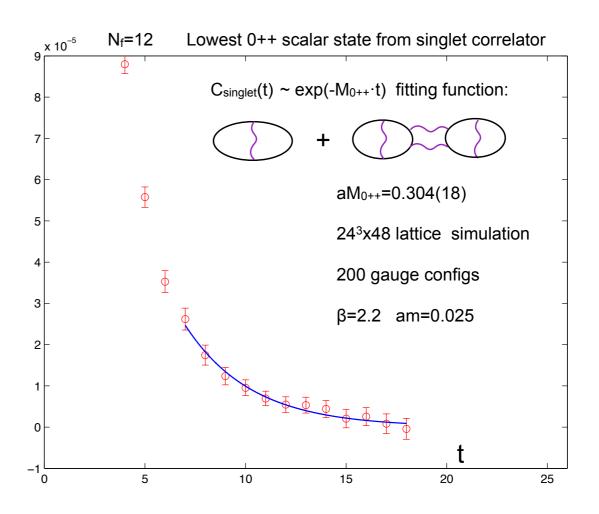
FL = 0.4 squeezed L= 0.8 fm, begins to look conformal not OK, misidentifies infinite volume phase

## The light 0++ scalar

### SCGT 2013-2015

### test of scalar technology:

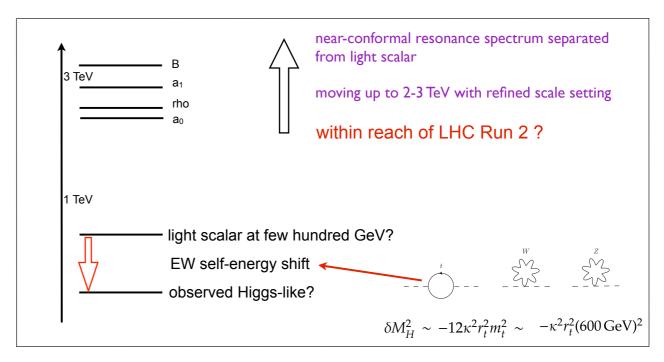




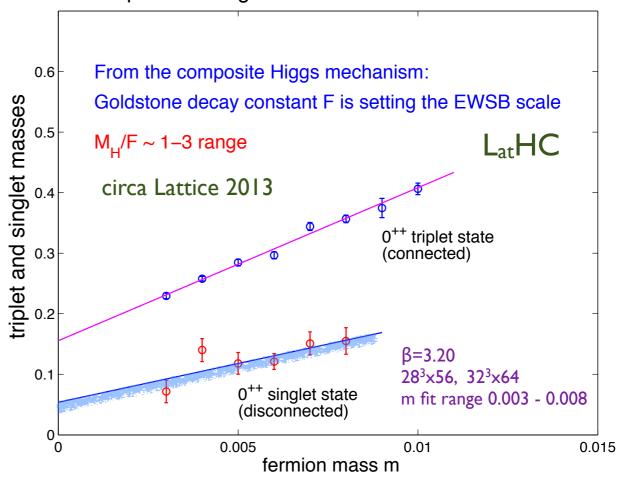
$$C(t) = \sum \left[ A_n e^{-m_n(\Gamma_S \otimes \Gamma_T)t} + (-1)^t B_n e^{-m_n(\gamma_4 \gamma_5 \Gamma_S \otimes \gamma_4 \gamma_5 \Gamma_T)t} \right]$$

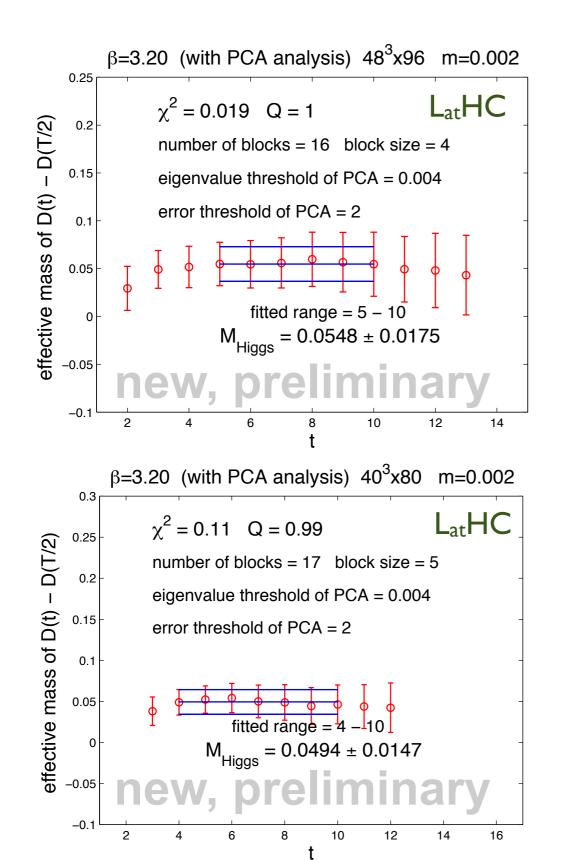
staggered correlator

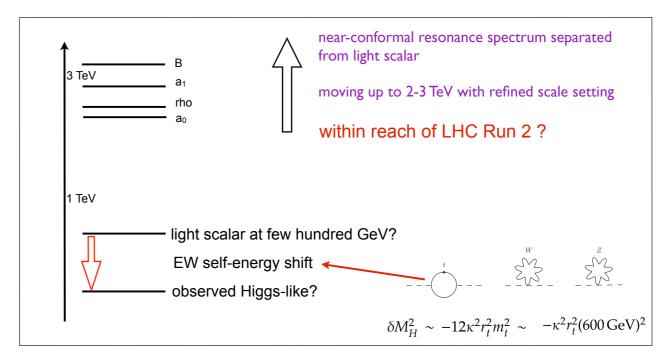
new results in  $N_f=2$  sextet model (this talk) and  $N_f=4/8/12$  models ( $L_{at}KMI$  talks, A. Hasenfratz talk)



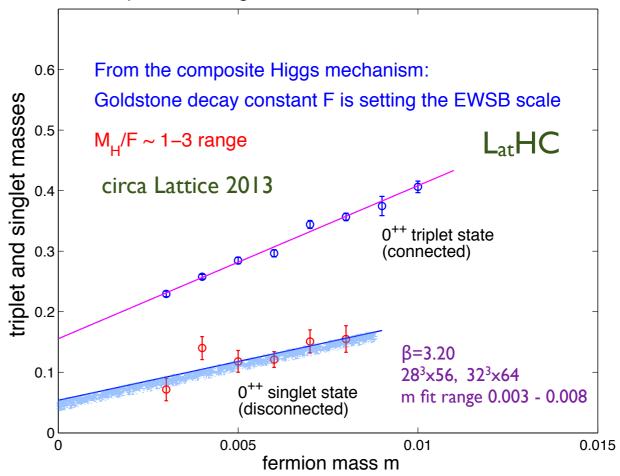
Triplet and singlet masses from 0<sup>++</sup> correlators

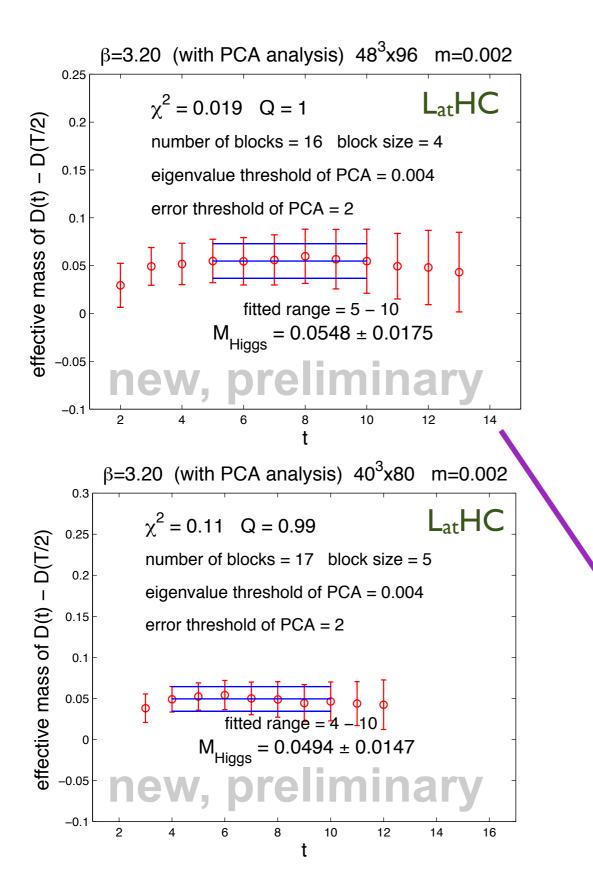


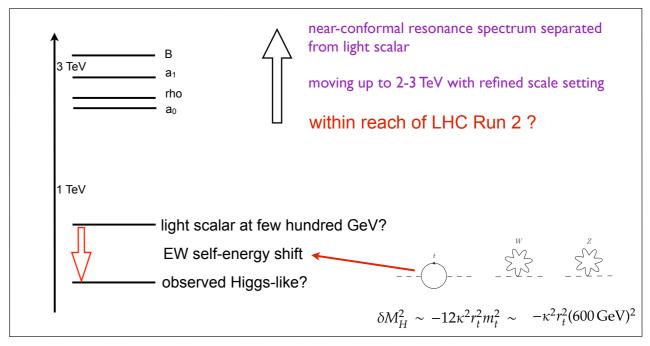




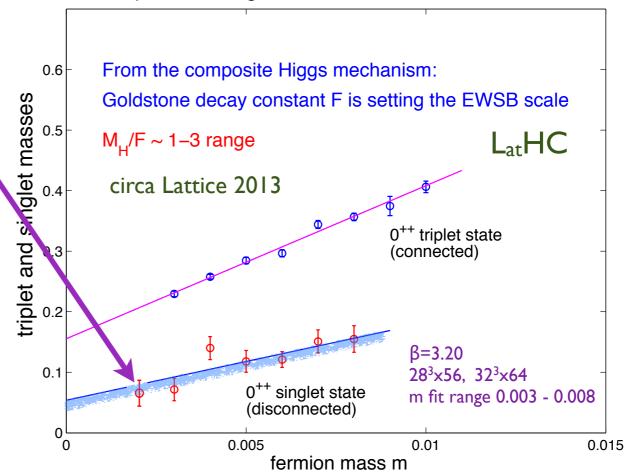


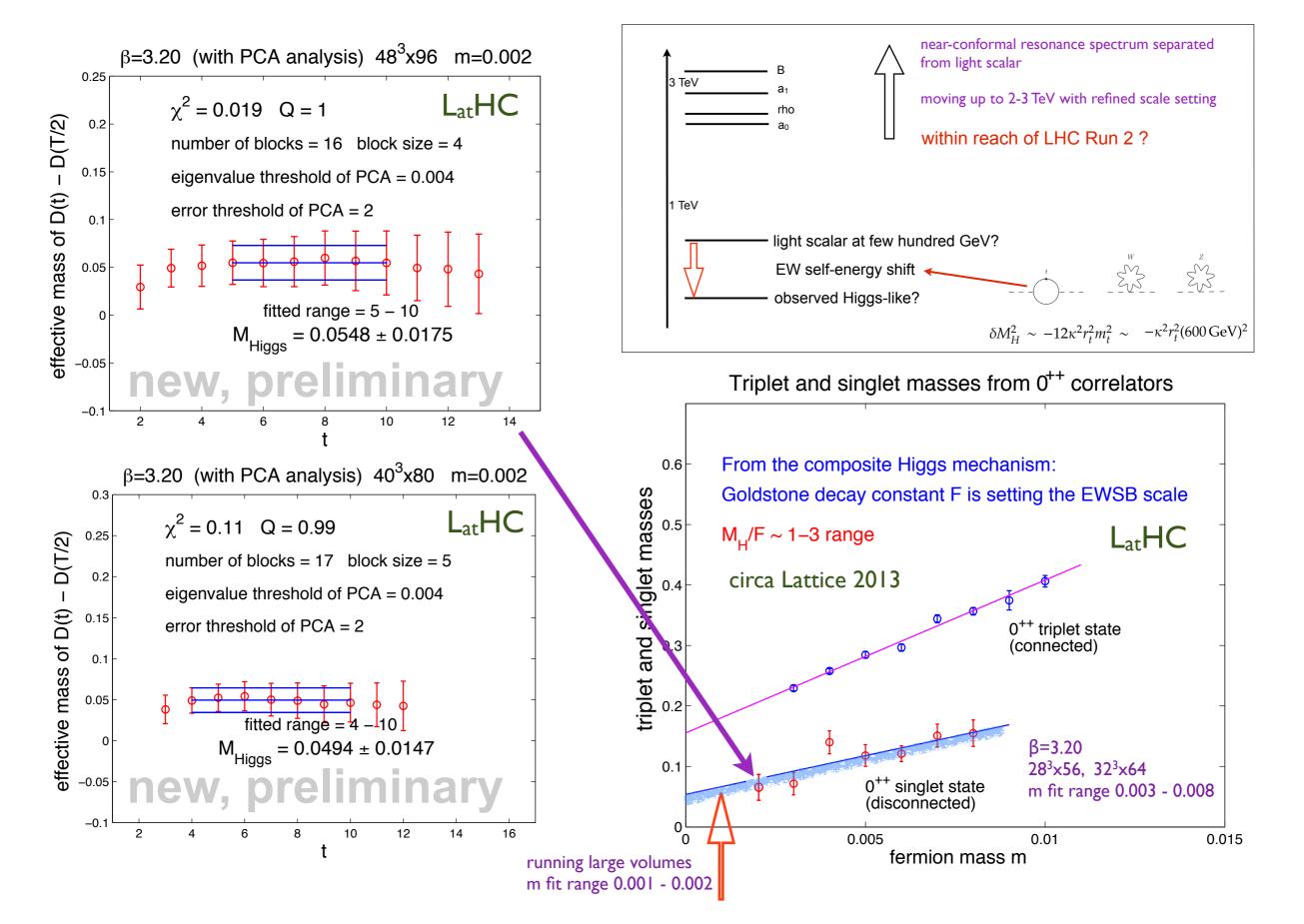


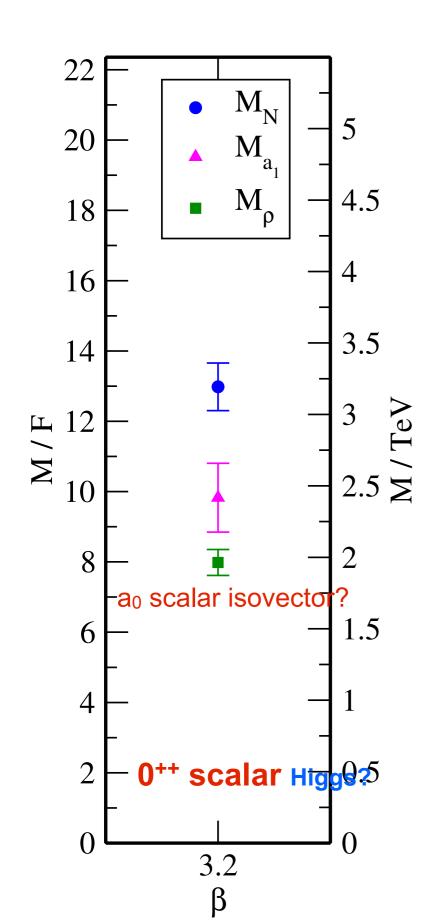




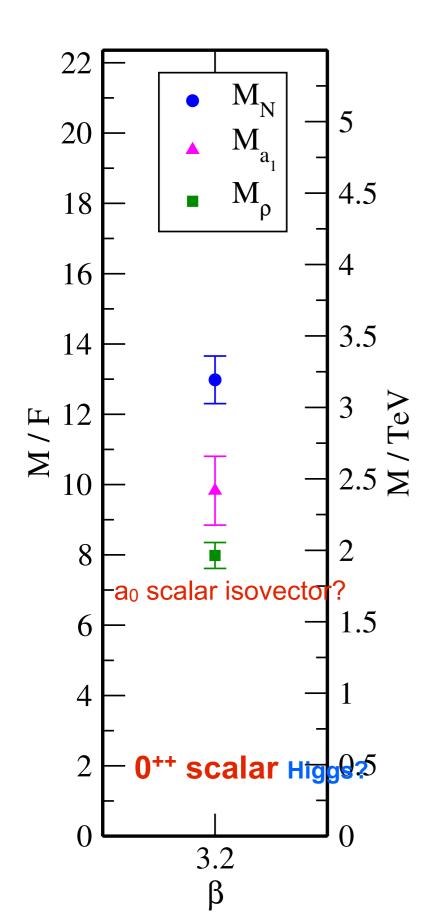
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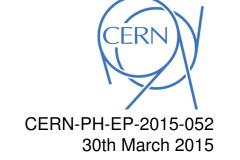


EUROPEAN ORGANISATION FOR NUCLEAR RESEARCH (CERN)





Submitted to: Eur. Phys. J. C.



Search for a new resonance decaying to a W or Z boson and a Higgs boson in the  $\ell\ell/\ell\nu/\nu\nu+b\bar{b}$  final states with the ATLAS Detector

The ATLAS Collaboration

with Sannino

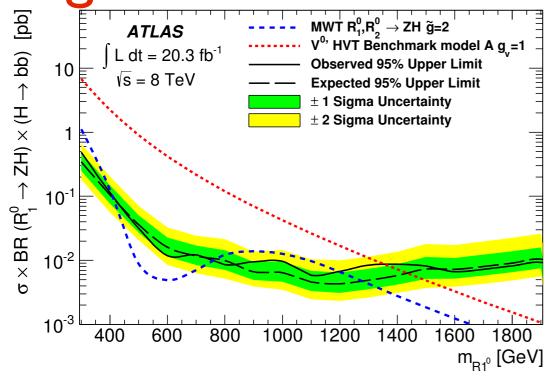
#### **Abstract**

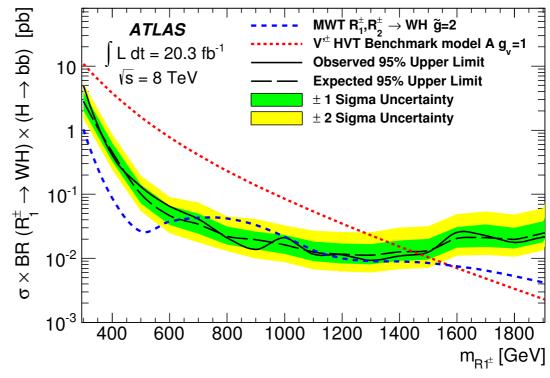
A search for a new resonance decaying to a W or Z boson and a Higgs boson in the  $\ell\ell/\ell\nu/\nu\nu+$   $b\bar{b}$  final states is performed using 20.3 fb<sup>-1</sup> of pp collision data recorded at  $\sqrt{s}=8$  TeV with the ATLAS detector at the Large Hadron Collider. The search is conducted by examining the WH/ZH invariant mass distribution for a localized excess. No significant deviation from the Standard Model background prediction is observed. The results are interpreted in terms of constraints on the Minimal Walking Technicolor model and on a simplified approach based on a phenomenological Lagrangian of Heavy Vector Triplets.

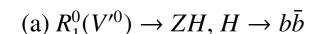
light 0++ scalar and spectrum

sextet model LatHC

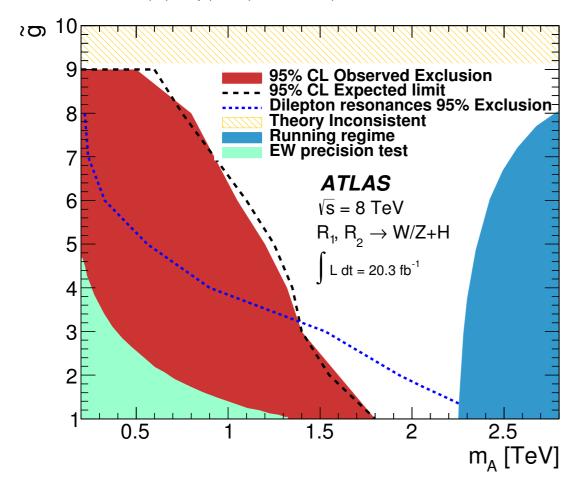


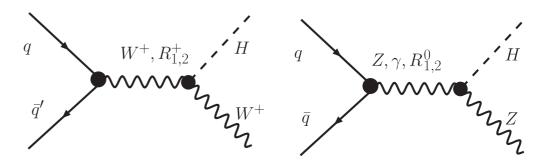






(b) 
$$R_1^{\pm}(V'^{\pm}) \rightarrow WH, H \rightarrow b\bar{b}$$

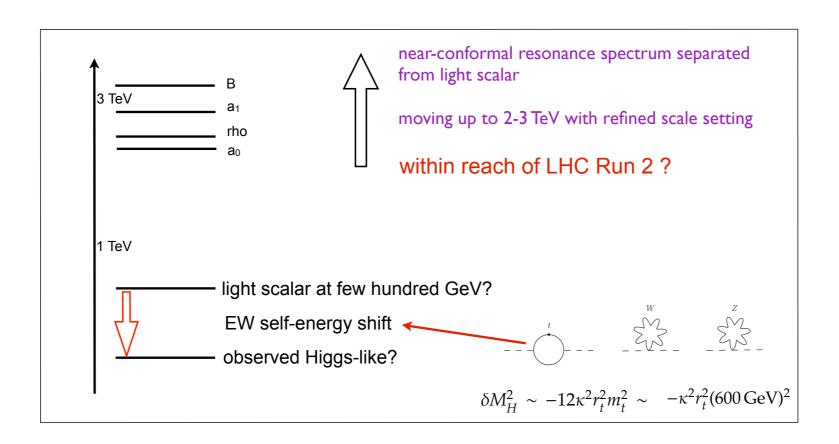




R<sub>1</sub> and R<sub>2</sub> couplings:

ĝ is the coupling in SU(4) vector boson g/g is the coupling to fermions

## light 0++ scalar and spectrum two tracks of challenges



### Theory track:

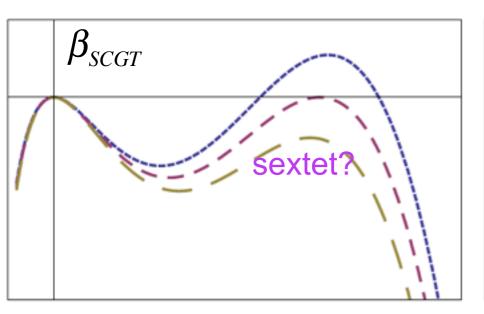
- is there a natural explanation for scale separation close to CW?
- testable meaning to dilaton interpretation?
- how to do mass deformed XPT when scalar is not decoupled from Goldstones?
- how the low mass scalar is effecting the RMT analysis in  $m \rightarrow 0$  limit ?

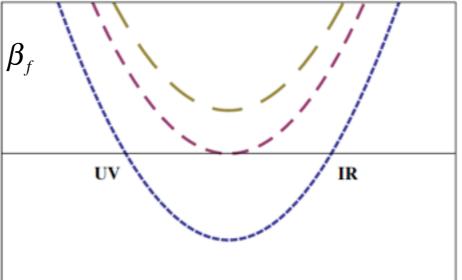
#### Simulation track:

- new mixed action strategy
- more accurate scale setting in continuum limit FL > !!
- analysis of slowly changing topology
- glueball mixing
- to reach decoupling of low mass scalar in RMT limit?

## an idea to work on?

BKT (Miransky) conformal phase transition?





tunable deformation of IRFP?

four-fermion operator with large anomalous dim?

$$L_{SCGT} \Rightarrow L_{SCGT} + \frac{f}{\Lambda^2} (\bar{\psi}\psi)^2$$

In the ivory tower we tune  $x = N_f / N_c$  in and out of CW starting from  $L_{SCGT}$  at IRFP and adding NJL term.

If anomalous dimension of  $(\overline{\psi}\psi)^2$  becomes marginal, the beta function  $\beta(g^2, f)$  can lead to collapse of the pair of the IR FP and the UV FP (created by the NJL term)  $\Rightarrow$  asymptotic safety. Only if **x** is tuned to **x**<sub>c</sub> critical of the BKT (conformal) phase transition. Miransky, Yamawaki Kaplan, Son, Stephanov Gies,... RG flow large-N double trace limit (Witten, Rastelli, Vecchi) Kutasov, ... (holographic)

NJL is misinterpreted but the general idea is attractive, does not need NJL:

Four-fermion interaction near four dimensions

J. Zinn-Justin \*

THE EQUIVALENCE OF THE TOP QUARK CONDENSATE AND THE ELEMENTARY HIGGS FIELD

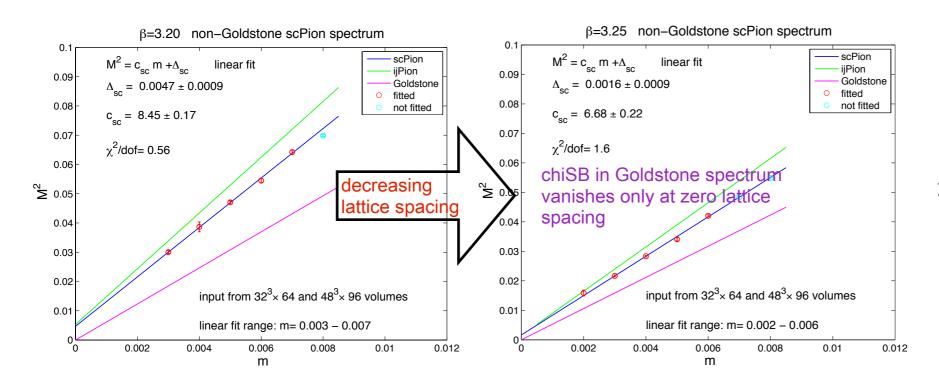
Anna HASENFRATZ

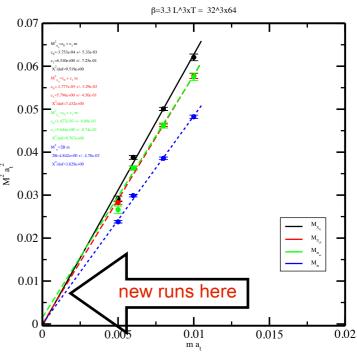
University of Arizona at Tucson, Department of Physics, Tucson, AZ 85721, USA

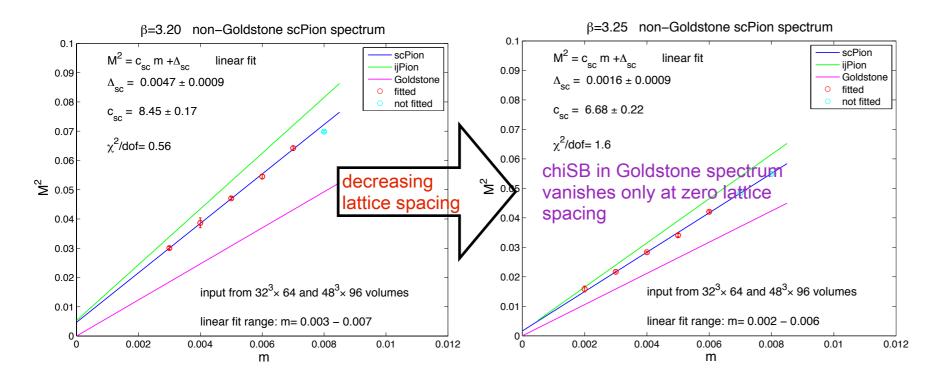
Peter HASENFRATZ\*, Karl JANSEN, Julius KUTI and Yue SHEN\*\*

On the lattice all terms are present on the cutoff scale in the Wilsonian sense and the model will decide what it wants to do with them.

Depending on anomalous dimension of  $(\bar{\psi}\psi)^2$  any of the scenarios can play out at any given point in the SCGT theory space.

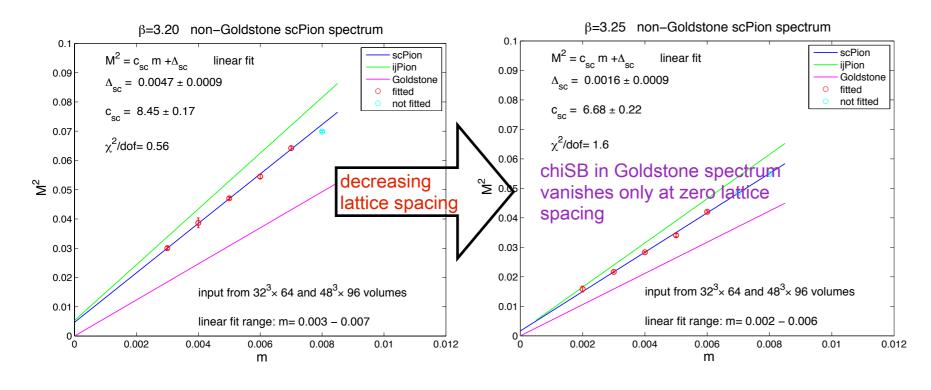


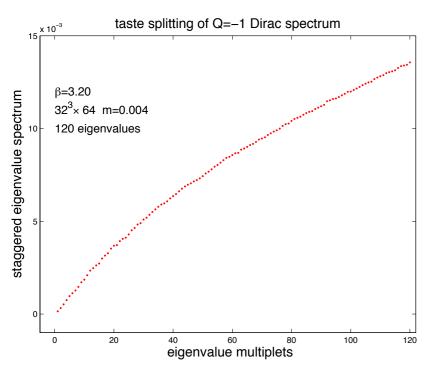


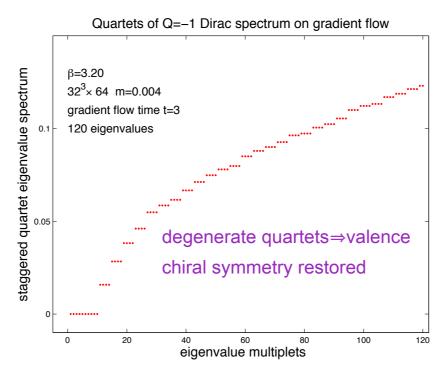


#### idea:

- use the gauge configurations generated with sea fermions
- taste breaking makes chiPT analysis complicated
- in the analysis use valence Dirac operator with gauge links on the gradient flow
- taste symmetry is restored in valence spectrum
- Mixed Action analysis should agree with original standard analysis when cutoff is removed: this is OK!

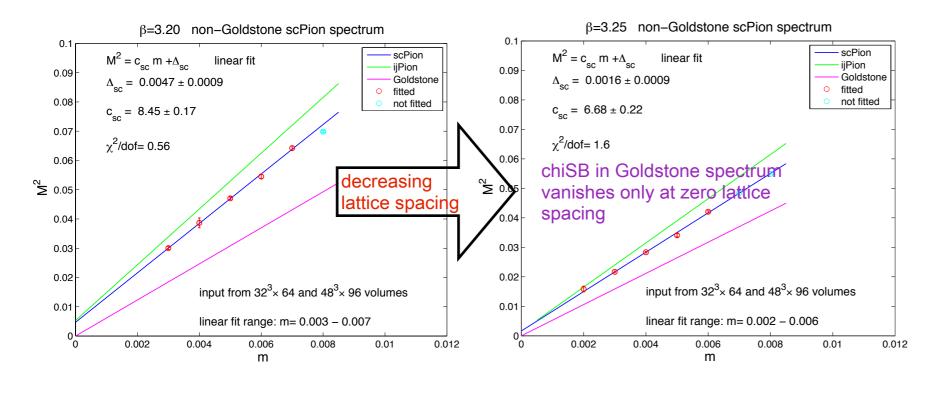






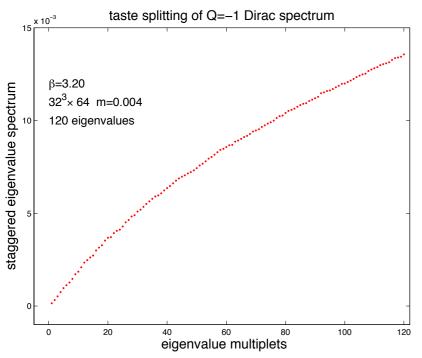
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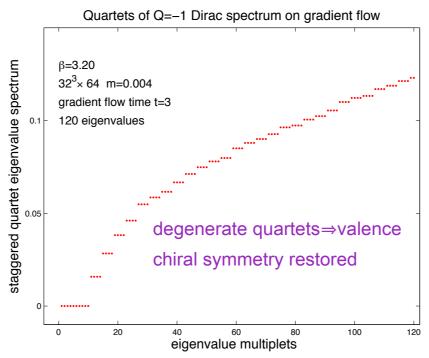
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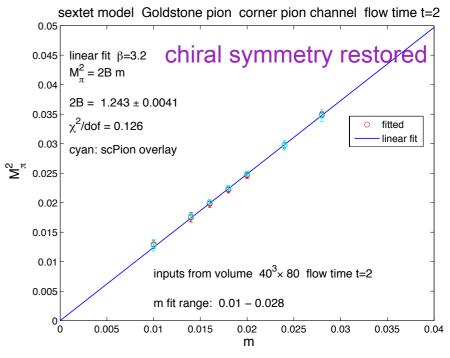


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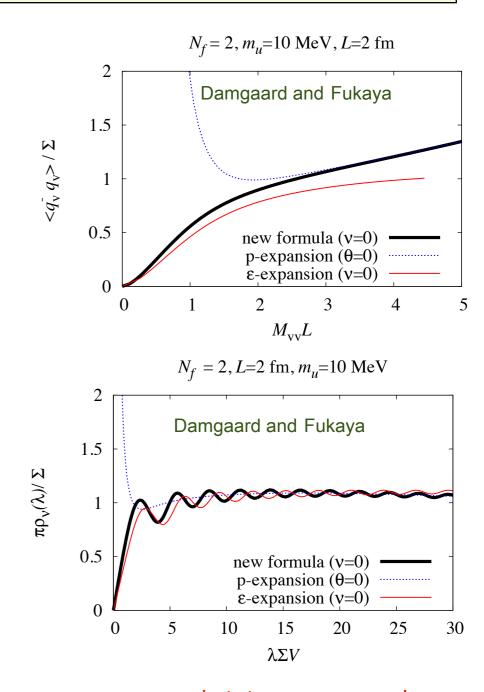
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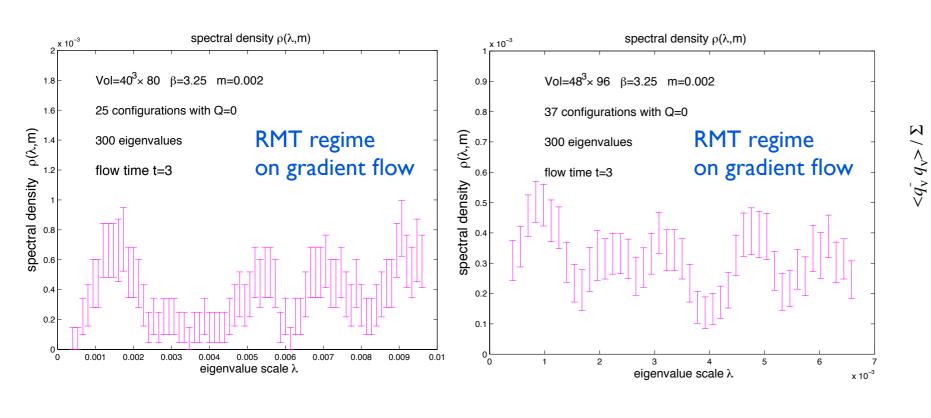


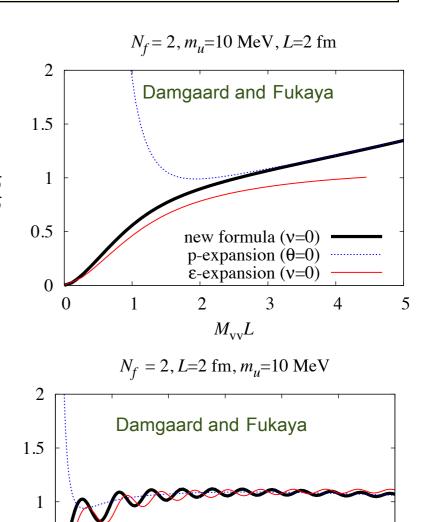
epsilon regime, p regime to epsilon regime crossover, valence pqChiPT with Mixed Action:



new analysis in crossover and RMT regime opens up with mixed action on gradient flow

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10

new formula (v=0) p-expansion ( $\theta$ =0)  $\epsilon$ -expansion (v=0)

20

25

15

 $\lambda \Sigma V$ 

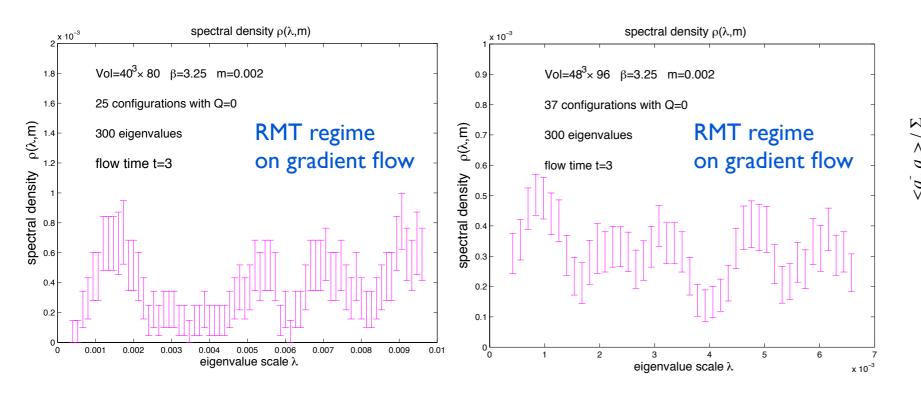
30

 $\pi \rho_{\nu}(\lambda)/\; \Sigma$ 

0.5

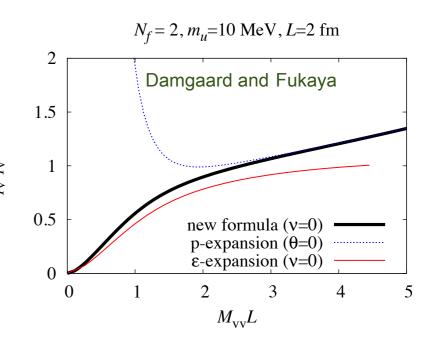
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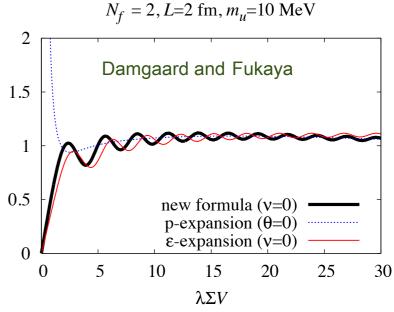
epsilon regime, p regime to epsilon regime crossover, valence pqChiPT with Mixed Action:





- GMOR implies large drop of order O(10) in the chiral condensate  $\Sigma$  is not RG invariant, requires renormalization
- in original analysis m  $\Sigma V \sim O(100\text{-}200)$  to reach RMT regime close to CW would require enormous resources
- in Mixed Action analysis  $\lambda \Sigma V \sim O(10-20)$  RMT regime can be reached





 $\pi \rho_{\nu}(\lambda)/\; \Sigma$ 

new analysis in crossover and RMT regime opens up with mixed action on gradient flow

#### The chiral condensate new method

chiral condensate and RG:

mode number distribution of Dirac spectrum

$$\rho(\lambda,m) = \frac{1}{V} \sum_{k=1}^{\infty} \left\langle \delta(\lambda - \lambda_k) \right\rangle \qquad \qquad \lim_{\lambda \to 0} \lim_{m \to 0} \lim_{V \to \infty} \rho(\lambda,m) = \frac{\Sigma}{\pi} \qquad \qquad \text{spectral density}$$
 (Banks-Casher)

$$\lim_{\lambda \to 0} \lim_{m \to 0} \lim_{V \to \infty} \rho(\lambda, m) = \frac{\Sigma}{\pi}$$

$$\nu(M,m) = V \int_{-\Lambda}^{\Lambda} \mathrm{d}\lambda \, \rho(\lambda,m), \qquad \Lambda = \sqrt{M^2 - m^2} \qquad \text{mode number function}$$

$$\Lambda = \sqrt{M^2 - m^2}$$

$$\nu_{\rm R}(M_{\rm R}, m_{\rm R}) = \nu(M, m)$$

renormalized and RG invariant

(Giusti and Luscher)

spectral density  $\rho(t)$  from ensemble averages over the D<sup>†</sup>D matrix with dimension N

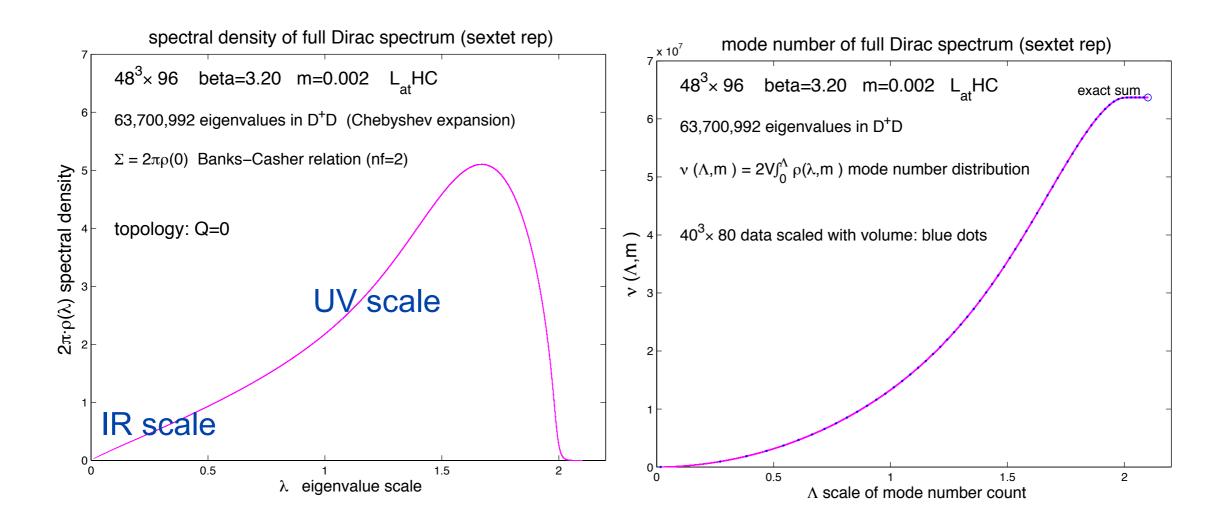
$$\rho(t) = \left\langle \frac{1}{N} \sum_{i=1}^{N} \delta(t - \lambda_i) \right\rangle_{\substack{\text{gauge} \\ \text{ensemble}}}$$

 $\rho(t) = \frac{1}{\sqrt{1-t^2}} \sum_{k=0}^{\infty} c_k T_k(t)$  expansion in Cebyshev polynomials

$$c_{k} = \begin{cases} \frac{2}{\pi} \int_{-1}^{1} T_{k}(t) \rho(t) & k = 0 \\ \frac{1}{\pi} \int_{-1}^{1} T_{k}(t) \rho(t) & k \neq 0 \end{cases} \implies c_{k} = \begin{cases} \frac{2}{N\pi} \sum_{i=1}^{N} T_{k}(\lambda^{2}_{i}) & k = 0 \\ \frac{1}{N\pi} \sum_{i=1}^{N} T_{k}(\lambda^{2}_{i}) & k \neq 0 \end{cases}$$

 $\sum_{k=1}^{N} T_k(\lambda^2)$  is given by trace of  $T_k(D^+D)$  operator

## The chiral condensate full spectrum



- nf=2 sextet example illustrates results from the Chebyshev expansion
- full spectrum with 6,000 Chebyshev polynomials in the expansion
- the integrated spectral density counts the sum of all eigenmodes correctly
- Jackknife errors are so small that they are not visible in the plots.

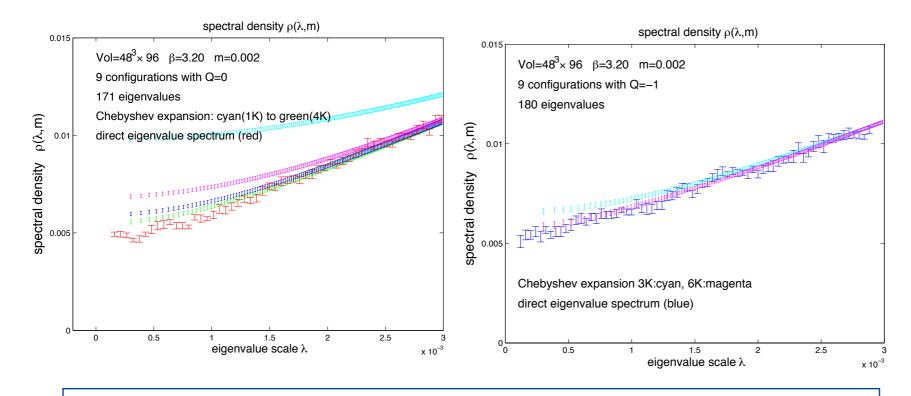
## The chiral condensate GMOR test in far IR

GMOR relation (nf=2):  $2BF^2 = \Sigma$  ( $\Sigma$  is the chiral condensate)

F: decay constant of Goldstone pion  $M_{\pi}^2 = 2B \cdot m$  in LO  $\chi$ PT

from chiral perturbation theory of the condensate in the p-regime:

$$\frac{\Sigma_{\rm eff}}{\Sigma} = 1 + \frac{\Sigma}{32\pi^3 N_F F^4} \left[ 2N_F^2 |\Lambda| \arctan \frac{|\Lambda|}{m} - 4\pi |\Lambda| - N_F^2 m \log \frac{\Lambda^2 + m^2}{\mu^2} - 4m \log \frac{|\Lambda|}{\mu} \right]$$



Improved determination of the chiral condensate  $\Sigma$  compared from Dirac spectra and the Chebyshev expansion.

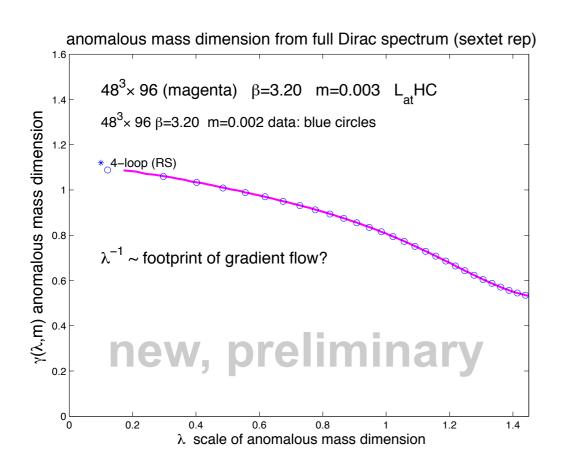
With the additive NLO cutoff term separated from B and new fit to F, the improved result on  $\Sigma$  eliminates previous discrepancies in the GMOR relation.

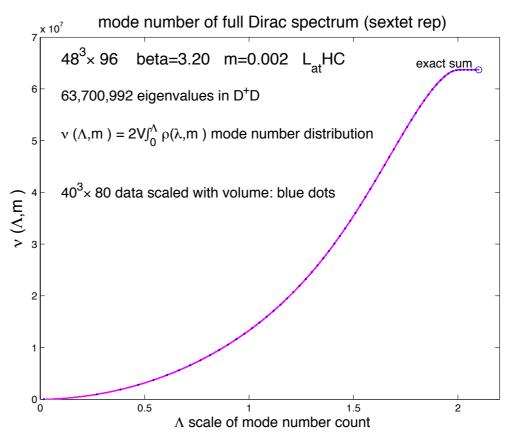
Boulder group pioneered fitting procedure

$$v_R(M_R, m_R) = v(M, m) \approx const \cdot M^{\frac{4}{1+\gamma_m(M)}},$$
or equivalently,  $v(M, m) \approx const \cdot \lambda^{\frac{4}{1+\gamma_m(\lambda)}}$ , with  $\gamma_m(\lambda)$  fitted

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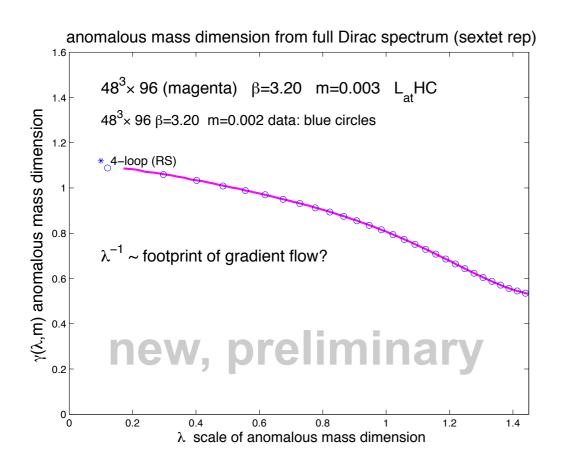
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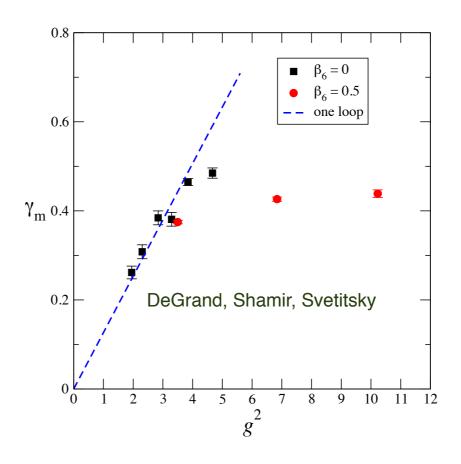




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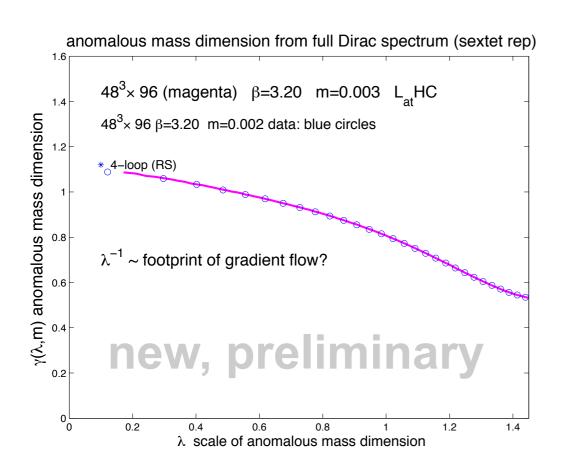


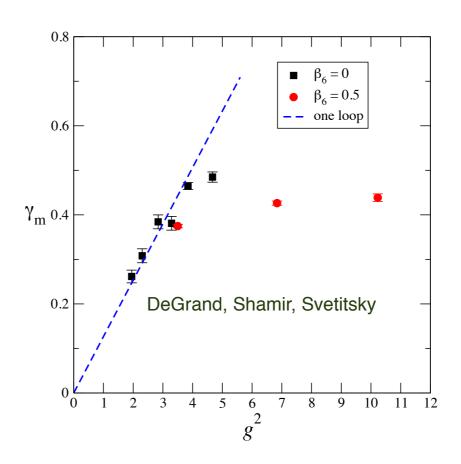


more details on the Nagoya poster!

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more details on the Nagoya poster!

How to match  $\lambda$  scale and  $g^2$ ?

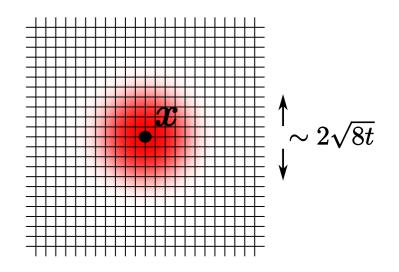
 $L_{at}HC$  group introduced the running coupling and its  $\beta$  function from the gauge field gradient flow with the scale set by the finite volume variations of it are becoming the standard approach

$$\dot{B}_{\mu} = D_{\nu}G_{\nu\mu} + \lambda D_{\mu}\partial_{\nu}B_{\nu}$$

$$B_{\mu,1}(t,x) = \int d^D y K_t(x-y) A_{\mu}(y),$$

$$K_t(z) = \int \frac{\mathrm{d}^D p}{(2\pi)^D} e^{ipz} e^{-tp^2} = \frac{e^{-z^2/4t}}{(4\pi t)^{D/2}}$$

Martin Lüscher earlier work by Neuberger



$$\langle E(t) \rangle = \frac{3}{4\pi t^2} \alpha(q) \{ 1 + k_1 \alpha(q) + O(\alpha^2) \}, \quad q = \frac{1}{\sqrt{8t}}, \quad k_1 = 1.0978 + 0.0075 \times N_f$$

t is the gradient flow time

Running coupling definition (range is  $(8t)^{1/2}$ ):

3rd Jacobi function

while holding 
$$c = (8t)^{1/2}/L$$
 fixed:  $\alpha_c(L) = \frac{4\pi}{3} \frac{\langle t^2 E(t) \rangle}{1 + \delta(c)}$   
$$\delta(c) = \vartheta_3^4(e^{-1/c^2}) - 1 - \frac{c^4 \pi^2}{3}$$

three different boundary conditions are used in practice:

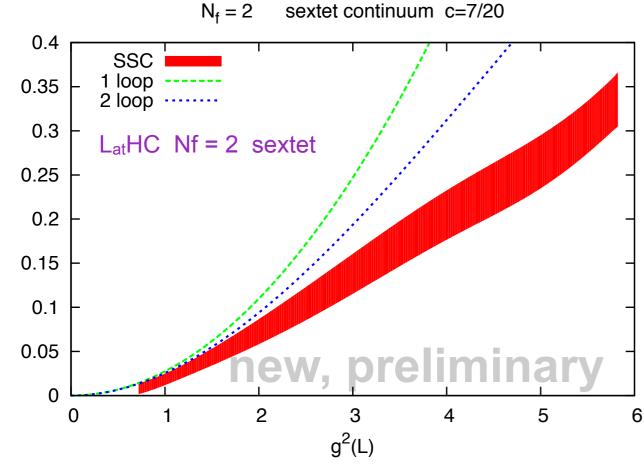
anti-periodic fermion fields Schrödinger functional twisted gauge fields and fermion fields

#### fundamental rep:

Nf=8 Boulder group and LatHC Nf=12 Boulder group and Taiwan group

sextet rep:

Nf=2 LatHC

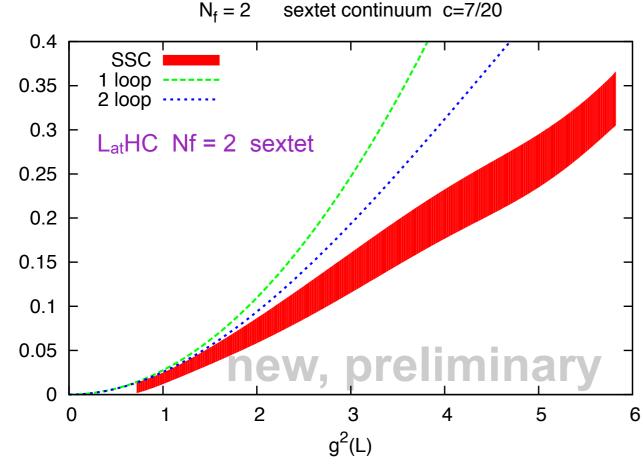


 $(g^2(sL) - g^2(L)) / \log(s^2)$ 

#### monotonic increase of beta function consistent with:

- mass deformed spectroscopy at low fermion mass
- chiral condensate
- GMOR
- mass anomalous dimension
- connection with g2(t,m) in bulk with chiSB

lattice step functions: 12→18, 16→24, 20→30, 24→36 last two step functions are critical in the analysis: SSC vs. WSC are consistent at large flow times which requires the large volumes



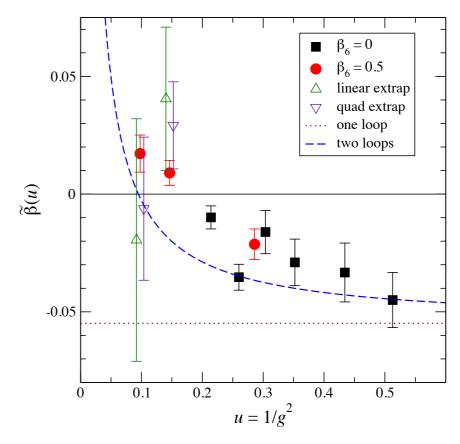
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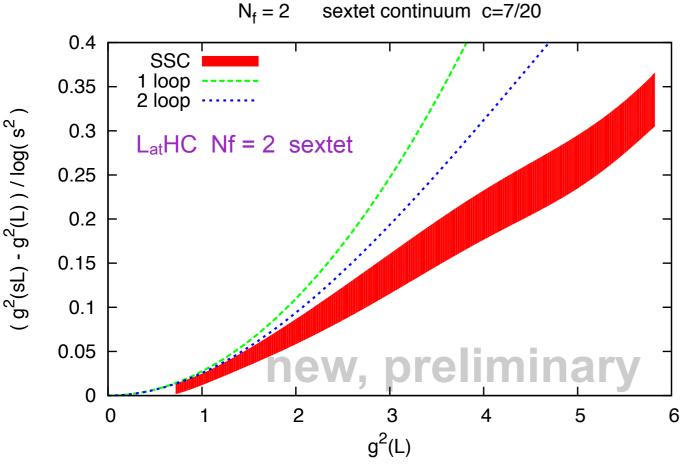
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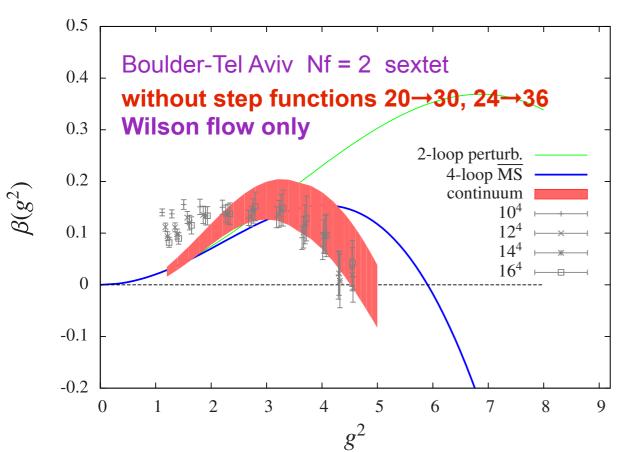


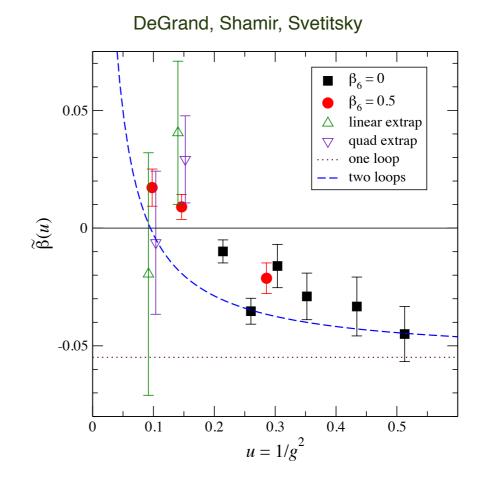


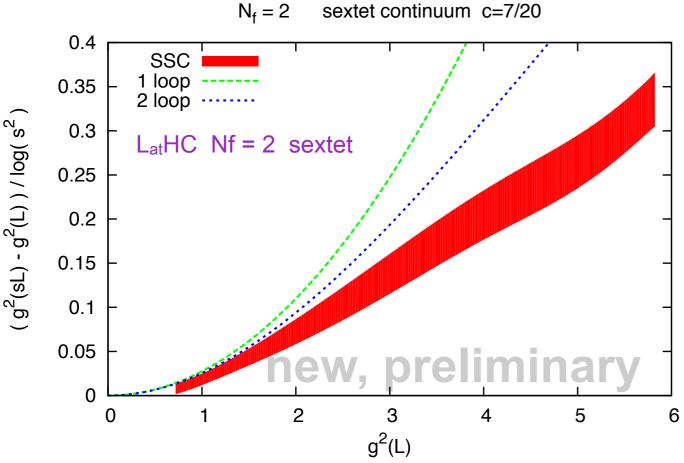
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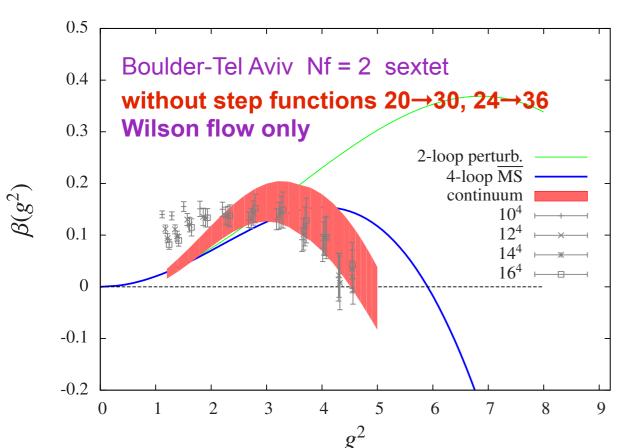




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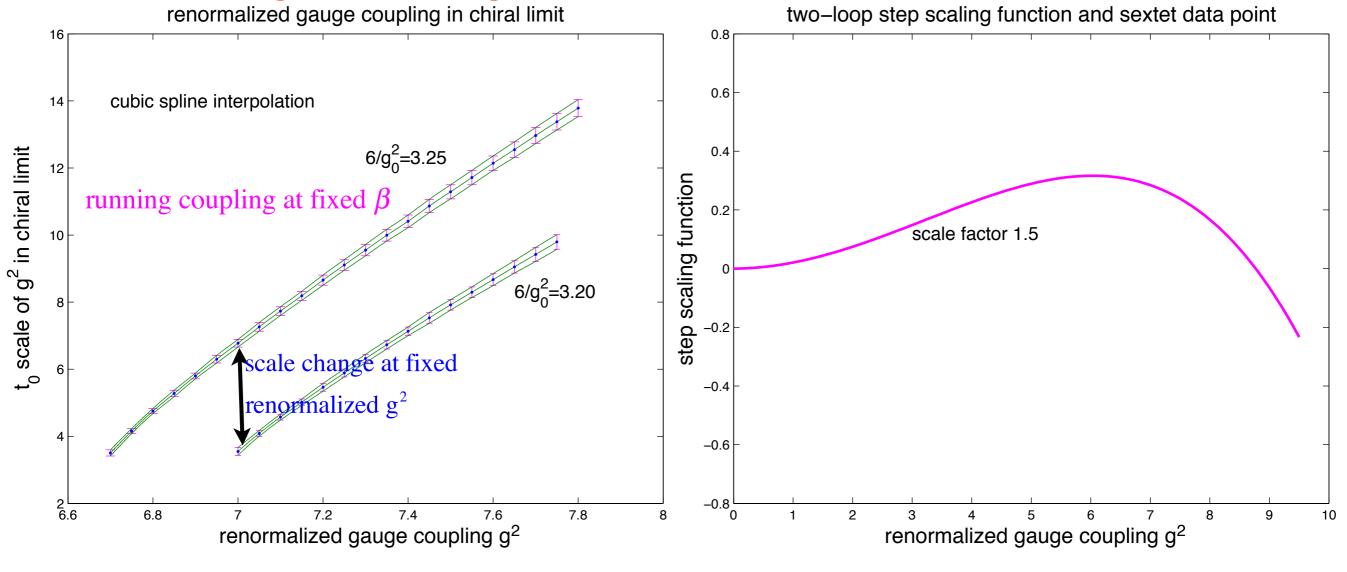
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This disagreement is between two groups and not about criticizing the staggered formulation

besides: promoting a beta function zero to conformal IRFP would require to remove the cutoff with the  $\boldsymbol{\omega}$  scaling exponent

· 11



running coupling, calculated at several bare  $g_0^2$ , allowes to determine the scale-dependent  $\beta$  function. This is in infinite volume, the opposite of running with a scale set by the finite volume.

leading dependence of  $g^2(t,m)$  on  $M_{\pi}^2$  is linear based on gradient flow chiPT Bär and Golterman works better than expected chiral logs are not detectable

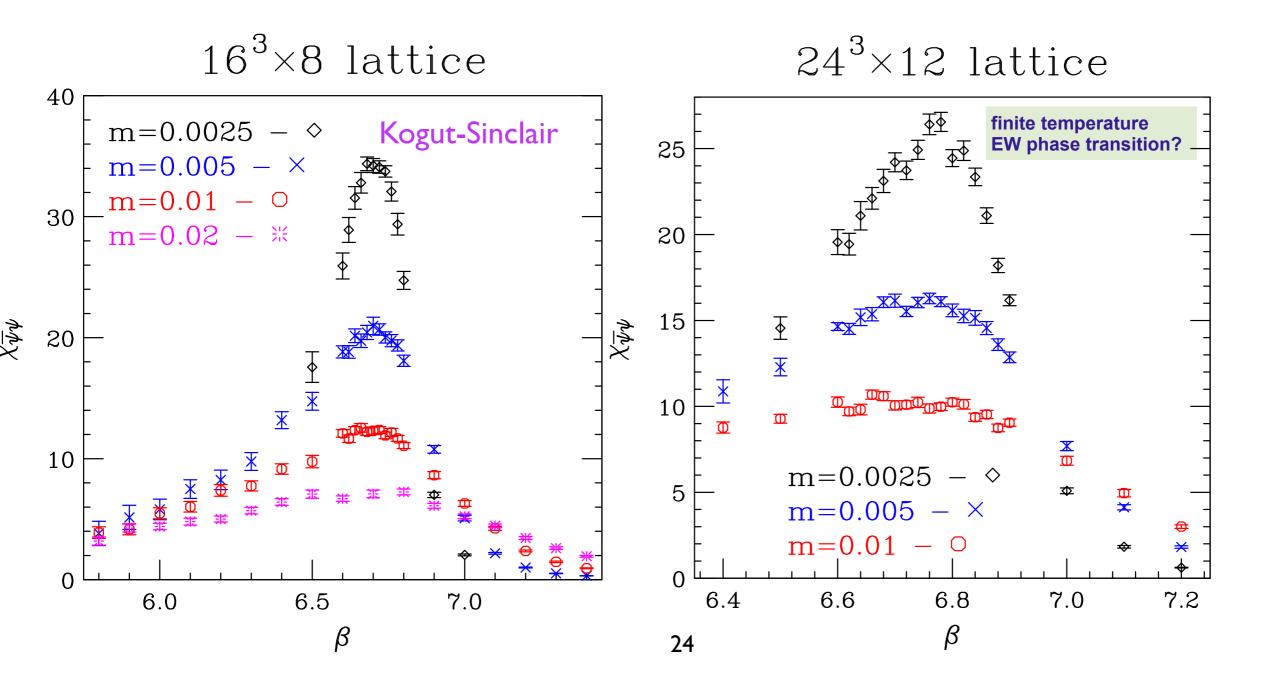
(not deep enough in leading log chiPT regime)  $\sim 2\sqrt{8t}$ 

## Early universe

Kogut-Sinclair work consistent with χSB phase transition

Relevance in early cosmology (order of the phase transition?)

LatHC is doing a new analysis using different methods



## Early universe

#### The Total Energy of the Universe:

Vacuum Energy (Dark Energy) ~ 67 %
Dark Matter ~ 29 %
Visible Baryonic Matter ~ 4 %

### Dark matter

self-interacting?

O(barn) cross section would be challenging

- lattice BSM phenomenology of dark matter
   Sannino and collaborators fundamental and adjoint rep
   LSD collaboration fundamental rep
- Nf=2 Qu=2/3 Qd = -1/3 fundamental rep udd neutral dark matter candidate
- dark matter candidate sextet Nf=2 electroweak active in the application
- 1/2 unit of electric charge (anomalies)
- rather subtle sextet baryon construction (symmetric in color)
- charged relics not expected?

Three SU(3) sextet fermions can give rise to a color singlet. The tensor product  $6 \otimes 6 \otimes 6$  can be decomposed into irreducible representations of SU(3) as,

$$6 \otimes 6 \otimes 6 = 1 \oplus 2 \times 8 \oplus 10 \oplus \overline{10} \oplus 3 \times 27 \oplus 28 \oplus 2 \times 35$$

where irreps are denoted by their dimensions and  $\overline{10}$  is the complex conjugate of 10.

Fermions in the 6-representation carry 2 indices,  $\psi_{ab}$ , and transform as

$$\psi_{aa'} \longrightarrow U_{ab} \ U_{a'b'} \ \psi_{bb'}$$

and the singlet can be constructed explicitly as

$$\mathcal{E}_{abc}$$
  $\mathcal{E}_{a'b'c'}$   $\psi_{aa'}$   $\psi_{bb'}$   $\psi_{cc'}$ .

### Summary: simplest composite scalar is probably very light (near conformality?)

light scalar (dilaton-like?) emerging close to conformal window?

spectroscopy
 emerging resonance spectrum ~ 2-3 TeV

• chiral condensate, large  $\gamma(\lambda)$  new method is very promising

running (walking) coupling in progress
 difficult, Gradient Flow is huge improvement

• Electroweak phase transition and baryon intriguing